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ANNUAL REPORT
ELECTRIC FIELDS IN EARTH ORBITAL SPACE

OCTOBER 1988

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*McDonnell Douglas Astronautics Company
Huntington Beach*

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Prepared by

W. P. Olson
K. A. Pfitzer

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→ southward and northward disturbances enter the magnetotail and then propagate within the magnetosphere. ~~We have developed~~ a magnetic field model that includes propagating waves within the magnetosphere. These models will help us develop a quantitative model of magnetospheric tail dynamics.

Keywords: → + fields developed

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Section 1

INTRODUCTION AND SUMMARY

Our magnetospheric work at MDAC over the past few years has centered on understanding physically the mechanisms for transfer of energy, mass, and momentum from the magnetosheath region into the magnetosphere, and on describing the interaction of the interplanetary magnetic field (IMF) with the magnetosphere. These are basic problems in magnetospheric physics since these processes must be understood before quantitative models of the many observed magnetospheric processes can be constructed. Such models are required for the prediction and specification of near earth orbital environmental "weather" parameters such as: trapped radiation fluxes, upper atmospheric density, ionospheric electron density, auroral particle precipitation and associated auroral luminosity, etc.

The question of particle entry into the magnetosphere has been one of the most fundamental problems in magnetospheric physics since it became an issue in the early 1960s. During the 1970s this question became even more basic to our understanding of the magnetosphere since there was growing correlative evidence that the magnetosphere was under the control of the solar wind and interplanetary magnetic field. The magnetospheric substorm was (and continues to be) of particular interest since the basic substorm question concerns whether or not the tail of the magnetosphere is controlled by the solar wind and interplanetary magnetic field or whether the plasmas in the tail can store energy and suddenly release it (a substorm) without external influences.

We have developed quantitative models of particle entry and solar wind and IMF influence on the magnetosphere. During this effort for ONR, our attention has been focused on the primary question: Is the substorm controlled by forces external to the magnetosphere or is it primarily a manifestation of processes occurring within the magnetospheric tail region?

Our models are driven by solar wind and IMF parameters. All parts of the models have been developed from first principles. It is our hope that in addition to shedding light on the substorm problem, the models will become useful in the truly predictive sense: that they can be used to anticipate the behavior of several of the routinely observed

magnetospheric features and their subsequent effects on the upper atmosphere and ionosphere and the orbital systems that operate in that region.

The solar wind "entry" problem and the substorm have both been important problems in magnetospheric physics since the 1960s. Currently, both problems are explained primarily in terms of "reconnection theory", which we believe has major flaws. We discuss the entry problem first.

1.1 Solar Wind Entry into the Magnetosphere

The first magnetospheric problem to be considered theoretically was the calculation of the size and shape of the magnetosphere. They were determined by equating the solar wind pressure with the energy density (pressure) of the magnetospheric magnetic field. To make this three dimensional problem tractable, an assumption was made; that the solar wind particles were all specularly (mirror like) reflected off of the geomagnetic field, \underline{B} . This approximation to the real field was quite reasonable since the interaction region of the incident particles with \underline{B} is very small (10s to 100s of km) compared to the scale size of the magnetospheric magnetic field (10s of thousands of km). This "pressure balance" formalism was used by many investigators to successfully determine the shape and size of the magnetosphere. The unfortunate legacy of this work, however, was the assumption that all of the solar wind particles are reflected off of the geomagnetic field and that none enters the magnetosphere.

Therefore, as early as 1960, concepts were already being suggested for the transfer of momentum across the boundary between interplanetary space (or more precisely, the magnetosheath) and the magnetosphere (the "viscous interaction" theories). Later, as the influence of the IMF on the magnetosphere became well established, the early work of Dungey and others on "reconnection" was suggested as another means for providing energy and mass to the magnetosphere from the solar wind. Several other concepts have been suggested: plasma instabilities along the magnetopause, a structured or "gusty" solar wind, and various diffusion processes.

All this time the assumption that solar wind particles were specularly reflected at the magnetopause was left basically unchallenged. Yet, it was well known that the higher energy solar cosmic rays readily gain access to the magnetosphere. Their entry was explained (correctly) by suggesting that since such particles have relatively large gyroradii, their interaction region with \underline{B} is large enough for them to sample the structure

of \underline{B} and therefore possibly gain access to the magnetosphere. (Recall that the path of a charged particle moving in a uniform magnetic field is circular. This fact forms the basis for the specular reflection assumption.)

However, it is known that the strength of \underline{B} along the magnetopause varies from about 75 nT at the nose of the magnetosphere to less than 5 nT in the distant magnetotail. Thus there is a gradient (structure) in \underline{B} (although it is admittedly small over a distance comparable to the gyroradius of a solar wind proton). It was reasonable then to ask: What is the lowest energy particle to gain access to the magnetosphere because of the existence of this known (small) gradient in \underline{B} ?

We have examined that question in quantitative detail and found that as the energy of the incident particle decreases, the range of directions through which it can enter the magnetosphere also decreases but that even for particles with solar wind energies the solid angle of allowed directions of incidence is finite. Thus, we have concluded that: The assumption of specular reflection may reasonably be used in the determination of magnetospheric size and shape but that some solar wind particles routinely gain access to the magnetosphere.

This work has suggested to us that all the diffusion and plasma instability theories may not be necessary in order to explain the interaction of the magnetosphere and solar wind; that instead all that was necessary was a more realistic examination of the interaction of the solar wind particles with \underline{B} .

We have been encouraged by this work since qualitatively it suggests that the solar wind particles will enter the magnetosphere along the sides of the tail and in the dayside cusp regions (where they are required). This contrasts with reconnection theory which suggests that particle entry is primarily at the nose of the magnetopause and over the lobes of the tail. Thus the reconnection theory requires the presence of a complex electric field structure within the magnetosphere in order to allow the entering particles to drift across \underline{B} to the regions where they are observed.

1.2 Problems with the Reconnection Theory

Another basic problem in magnetospheric physics is the understanding of the interaction of the IMF with the magnetosphere. Currently, most of the magnetospheric physics community chooses to represent this interaction in terms of reconnection theory.

Basically, reconnection theory supposes that the two magnetic field sources (the IMF and \mathbf{B}) become "tied together" in the presence of a plasma. This process has the consequence of connecting the two magnetic fields such that field lines emanating from one are ultimately joined to the other source. An important by-product is the acceleration of charged particles in the "reconnection region". Proponents of reconnection theory suggest that it can explain the entry of the required solar wind plasma into the magnetosphere and also the relatively hot plasma observed in the tail of the magnetosphere.

We have several problems with this theory, however, and with support from ONR have developed an alternative physical description of the IMF and its interaction with the magnetosphere. Some of the problems we see in reconnection theory are as follows. Reconnection is dependent on the direction of the IMF which is observed to change, typically on the order of every few hours. Thus it is not clear to us how any reconnection process can drive any of the processes observed to persist at all times in the magnetosphere (e.g., the plasma sheet, the several magnetospheric currents, etc.). Its proponents state that reconnection takes place predominantly at the nose of the magnetopause and over the lobes of the tail. This poses the requirement of maintaining a complex electrostatic field at all times but in varying directions in order to permit particles entering there to journey to those regions where they are required (e.g., the plasma sheet). Finally, there is no conclusive evidence that this process operates in the IMF and magnetosphere even though a concerted effort has been expended searching for examples. A case in point was the AMPTE ion releases made upstream of the bow shock. No evidence was found for the presence of any ions entering near the nose of the magnetosphere even though the releases were performed during intervals of southward IMF (conditions favorable for dayside reconnection).

Section 2

RECENT ACTIVITIES

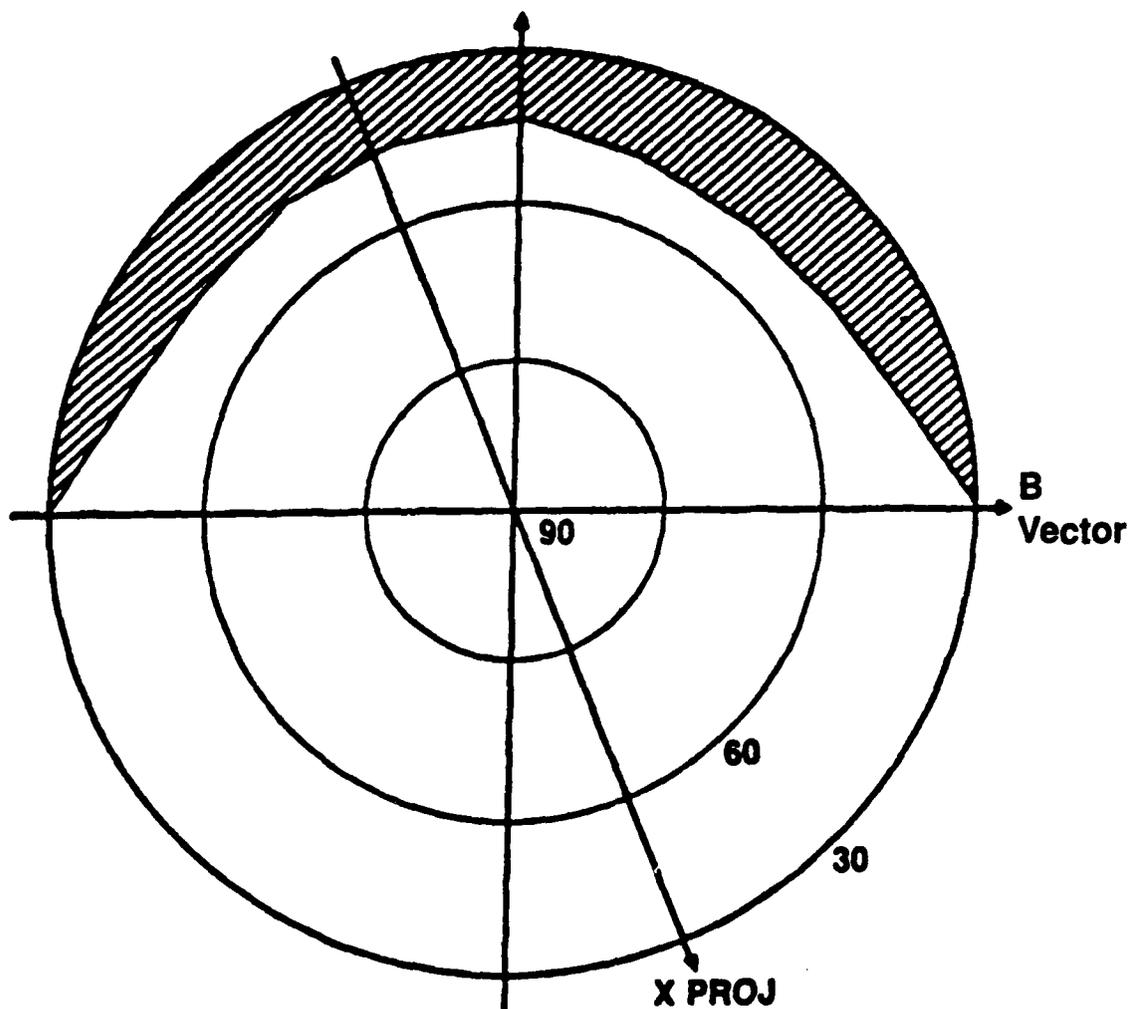
We have recently worked on the entry of solar wind into the magnetosphere and on the control the interplanetary magnetic field (IMF) has on the magnetosphere. Correlative data conclusively prove that both the solar wind and the IMF exert considerable influence on magnetospheric processes, and the magnetospheric substorm in particular.

Our work on this subject may be roughly divided into two parts, particle entry and representation of the interaction of the IMF with the magnetosphere. We discuss particle entry first.

2.1 Particle entry

We have quantitatively determined where on the magnetopause that low energy (solar wind) charged particles can gain entry to the magnetosphere. To do this, we used a realistic quantitative model of the magnetospheric magnetic field. Particles of the same energy but differing incidence angles were then introduced to this field at a point on the magnetopause. It was found that at most points a finite "entry cone" persists. The entry cone is defined as that region (represented as a solid angle) through which particles have access to the magnetosphere. The value of the entry cone varies with particle energy and location on the magnetopause. We found that the entry cone was largest along the equatorial flanks of the tail. (An example of the entry cone is shown in Figure 1 and the size of the entry cone over the flanks of the tail is shown for protons in Figure 2.) The entry cone is exactly zero along the intersection of the noon-midnight meridian with the magnetopause by symmetry since there the particles really are specularly reflected.

Study of entry cone size suggests that no magnetosheath particles enter the magnetosphere over the lobes of the tail and also that no particles enter at the nose of the magnetosphere. The regions where entry does readily occur are along the sides (flanks) of the tail and in the vicinity of the dayside cusps. The "gradient drift" entry mechanism therefore supplies solar wind plasma directly to those regions of the



Boundary Position: -20.00, -17.90, 3.16
1 keV Protons
Entry Cone = 0.249 Sr

Figure 1. Description of an entry cone. The shaded portion represents the range of impact directions over which the particle gains entry to the magnetosphere. The boundary location is down the tail just above the magnetic equatorial plane (along the dawn flank). The projection of the x axis (in solar magnetospheric coordinates) and the direction of \underline{B} at the entry point are shown. Particles moving away from the sun with a near grazing incidence are shown as "allowed" impact angles.

magnetosphere where plasmas are observed (e.g., the plasma sheet in the tail and the dayside cusps). This is unlike the reconnection theories which introduce plasma near the nose of the magnetosphere and over the lobes of the tail. To date this work on particle entry has shown that particle entry depends on the strength of \underline{B} , its structure, and on the particle distribution function.

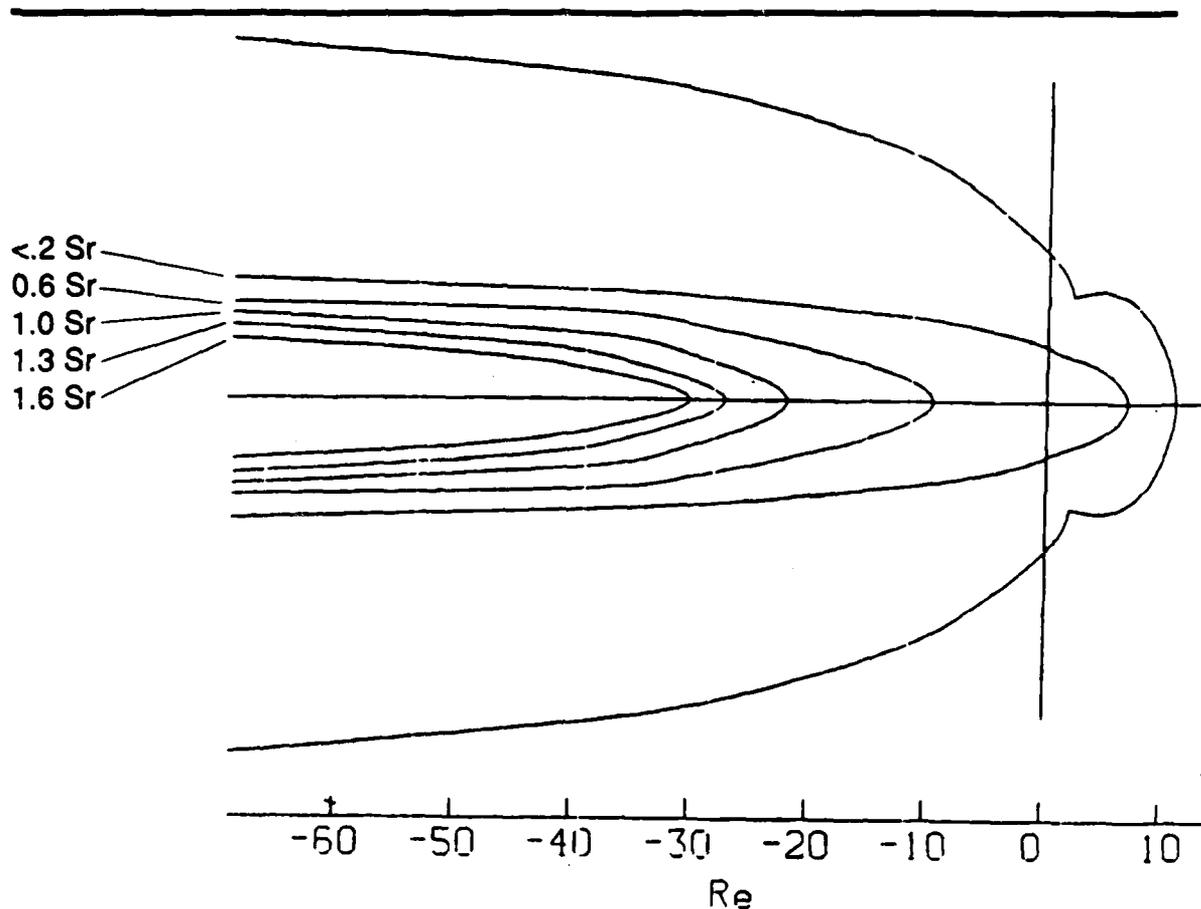


Figure 2. Contours of constant entry cone magnitude on the dawn side of the magnetosphere.

2.2 Representation of the Interplanetary Magnetic Field (IMF)

Correlative data establish beyond doubt that the IMF exerts considerable influence on magnetospheric processes. This subject has been exhaustively studied in terms of reconnection theory. We, however, have several complaints with reconnection theory (it introduces particles into the wrong regions of the magnetosphere: the IMF [and therefore the reconnection process] is continuously changing direction; there is little or no direct magnetospheric observational evidence for its existence etc.); and have, therefore, attempted to explain the interaction of the IMF with the magnetosphere in a way that we can understand the physics.

Basically, we have represented the IMF as the superposition of electromagnetic waves. The only limitation of our investigation is that we must restrict it to wavelengths

that are small with respect to the scale size of the magnetosphere. Thus we can examine only waves (or disturbances) with periods less than one hour. Although limited in this regard, our study sheds light on many aspects of the interaction of the IMF with the magnetosphere. These are briefly reviewed. We first summarize our findings on the propagation of electromagnetic disturbances in the solar wind.

The interplanetary field (represented as the superposition of several periodic electromagnetic waves) can persist only in the presence of the solar wind plasma. In the absence of a plasma, the presence of a time varying magnetic field in interplanetary space of the magnitude of only a few nT would have associated with it an electric field with a magnitude on the order of 1 volt per meter, which is at least three orders of magnitude larger than the magnitude of electric fields observed in the interplanetary region.

In the solar wind, in the absence of a background ambient magnetic field, any disturbances with periods on the order of minutes to hours will be rapidly attenuated unless they are driven continuously in a local region. Electromagnetic disturbances in the solar wind can persist only in the presence of an ambient (lower frequency) magnetic field. When such conditions are present (the presence of both plasma and "background" magnetic field) electromagnetic waves with periods from a few minutes to several hours can propagate over distances large with respect to the magnetosphere size without appreciable attenuation. The background (ambient) magnetic field is provided by the "solar sector magnetic field" which is co-produced with the solar wind. It moves outward from the sun with the solar wind, and has a period of about two weeks - much longer than the characteristic periods of the electromagnetic disturbances being considered. Electromagnetic waves allowed in the interplanetary medium propagate at the Alfvén speed. There are two wave modes that propagate without appreciable attenuation.

1. When the propagation vector is parallel to the ambient magnetic field.
2. When both the propagation vector and the disturbance electric field are perpendicular to the ambient magnetic field direction.

Each of these cases is represented schematically in Figure 3. It then was a simple exercise to examine the interaction of such waves when they encountered a discontinuity in plasma (the magnetopause). At the magnetopause these electromagnetic

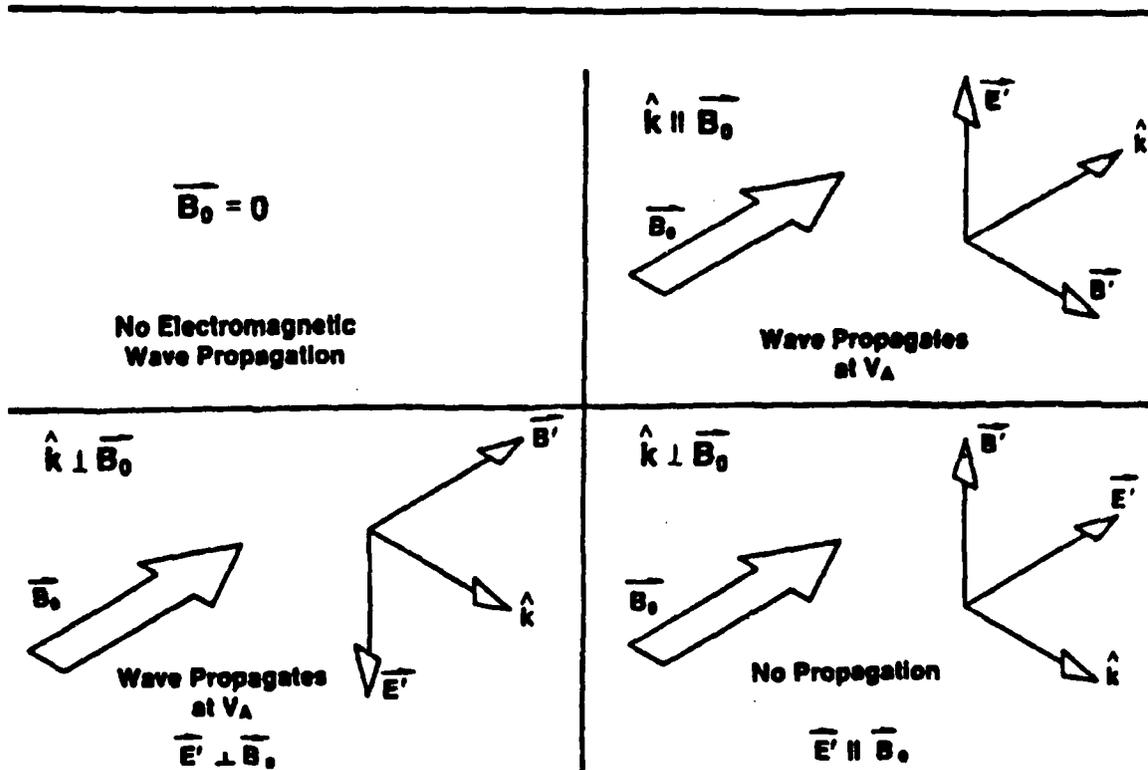


Figure 3. Electromagnetic wave propagation in a tenuous plasma. The wave propagates with little attenuation when its propagation vector is parallel to the direction of the ambient magnetic field, and when the propagation vector and the disturbance \vec{E} vector are perpendicular to \vec{B} .

disturbances are reflected and refracted. A portion of the field can penetrate the magnetosphere. The properties of the penetrating field are determined by magnetospheric parameters (e.g. plasma density, "ambient" magnetospheric magnetic field, etc.). Penetration of these interplanetary magnetic disturbances into the magnetosphere occurs most readily where the magnetic field strength is low and the plasma density relatively high. Thus, the flanks of the tail are the primary region for IMF penetration into the magnetosphere. An example of allowed transmission is shown in Figure 4 where B_0 represents the ambient fields in the magnetospheric and interplanetary regions and B' and E' define the disturbance. Note that the directions of B' and E' do not immediately change as the disturbance enters the magnetosphere and that E' is always perpendicular to B_0 .

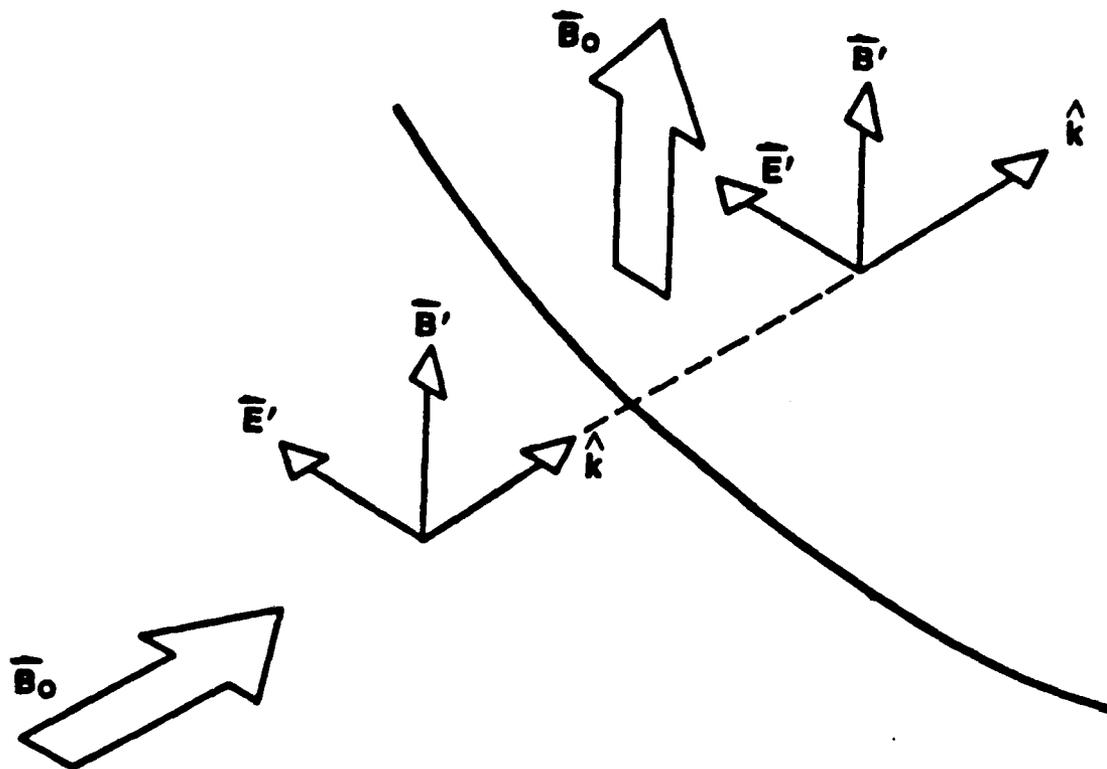


Figure 4. Transmission through the magnetopause and propagation in the magnetosphere. This process is completely analogous to the passage of light from vacuum into glass. In both media (the interplanetary region and in the magnetosphere) the disturbance \underline{E} vector must be perpendicular to the ambient magnetic field direction.

Magnetic disturbances in the north or south direction (perpendicular to the magnetic equatorial plane and parallel or antiparallel to the ambient magnetospheric field resident there) can most significantly influence the tail plasma sheet region. The entry of a northward disturbance field decreases the beta (ratio of particle kinetic energy to magnetic field energy density) while a southward field increases beta. We note that for a given ambient magnetospheric field, the entry of a southward disturbance field has more effect than that of a northward disturbance of the same magnitude on the percent change in beta. Thus this work on a physical representation of the IMF has already suggested why a southward directed IMF is observed to have such a large influence on the magnetospheric substorm process.

Section 3

SUMMARY OF WORK PERFORMED IN THE PAST YEAR

The basic problem of magnetospheric physics remains the explanation of the response of the magnetosphere to variations in the solar wind and the IMF. The work that has been completed as well as the work already in progress places us in a position to take a large step toward the goal of solving that problem. We now understand how the solar wind enters (and influences) the magnetosphere. We have also developed a physical description for the interaction of the IMF (represented as the superposition of a set of periodic electromagnetic disturbances) with the magnetosphere, and are studying the dynamics of energization processes in the magnetotail.

In our earlier work for ONR, we showed that variations in the IMF can be represented as a superposition of a set of periodic electromagnetic disturbances propagating in the magnetized interplanetary plasma, the solar wind. The solar wind is represented as a tenuous plasma with an imbedded magnetic field, the well known solar sector structure magnetic field. Higher frequency waves can propagate in this medium with the Alfvén speed if their propagation direction is parallel to the sector structure magnetic field. Such waves can also propagate if the propagation vector is perpendicular to the sector structure magnetic field providing the wave's electric field. The study also determined that when the wave interacts with the magnetopause, the wave is either reflected or transmitted. The study determined that transmission was possible only near the equatorial flanks of the tail and near the cusp regions, regions of weak magnetic field and where the plasma density inside the magnetopause was relatively high. Once the wave has penetrated through the boundary of the magnetosphere, it will either travel at the local Alfvén speed or be absorbed. It was determined that only waves whose magnetic field vector oscillates in the north/south direction are capable of propagating in the interplanetary medium, pass through the boundary near the equatorial flanks and then propagate within the magnetospheric tail field region.

Variations in the north south component of the interplanetary field have long been associated with substorm triggers. To begin the study of substorm trigger, studies of the magnetic field variations in the magnetospheric tail in response to changes in the IMF were initiated. The effort during this past year focused on the development of a

model of the magnetospheric tail that included an arbitrary north south wave propagating through the tail at the Alfvén speed. A newly revised dynamic magnetic field model was used as the starting point of the study.

The new dynamic magnetic field model consists of four separate magnetic field routines, an internal field routine, a magnetopause magnetic field routine, a ring current magnetic field routine, and a tail current magnetic field routine. The model developed for the IMF wave entering the magnetopause and then propagating within the magnetosphere was written as a time dependent routine which was then combined with the other four existing routines to produce a time dependent magnetic field model.

3.1 The IMF Pulse Within the Magnetosphere

In our previous work we showed that northward or southward changes in the IMF are able to enter the magnetosphere along the flanks of the tail. Once the disturbance pulse is within the magnetosphere it propagates down and across the tail at the Alfvén speed. The primary effort during this past year was the development of a dynamic model of the magnetosphere that included a propagating wave within the magnetosphere.

The model of the disturbance pulse developed during this past year includes the following options:

1. The IMF disturbance can either be north or south.
2. The rise time of the disturbance vector is adjustable. The shape of the rising edge of the disturbance vector was assumed to have a sinusoidal shape (quarter period) of arbitrary duration.
3. The duration of the pulse is adjustable. If the pulse has a finite length, the fall time of the pulse is adjustable. The falling portion of the pulse is also a quarter of a sinusoid.

A specified disturbance having a profile determined by the above constraints is initiated outside the magnetosphere and then allowed to interact with the magnetopause. The wave which enters the magnetosphere near the equator and along the flanks is then allowed to modify the magnetospheric magnetic field.

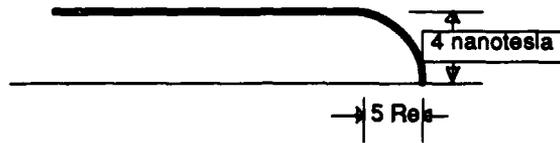
In this progress report we present four cases:

1. A 4 nanotesla northward disturbance with a rise time of $5 R_E$ and an infinite pulse length
2. A 4 nanotesla northward disturbance with a rise time of $5 R_E$ and a fall time of $5 R_E$
3. A 4 nanotesla southward disturbance with a rise time of $5 R_E$ and an infinite pulse length
4. A 4 nanotesla southward disturbance with a rise time of $5 R_E$ and a fall time of $5 R_E$

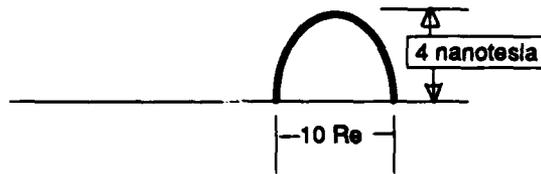
Figure 5 shows the space profile of these pulses. Since the pulse is traveling through the IMF with the Alfvén speed, the pulse has definite spatial extent. When the leading edge of the pulse is at a given location, the pulse maximum is some distance upstream. For presentation purposes it is easiest to display the figures as a function of location with each figure describing the magnetic field topology at a fixed time. The pulses used for these test cases were assumed to have their propagation vectors in the ecliptic plane and make a 45 degree angle with the sun-earth line. The effect of these disturbance pulses on the magnetospheric magnetic field is shown by displaying the behavior of the magnetic lines of force in the noon-midnight meridian plane.

3.1.1 Northward Pointing Disturbance

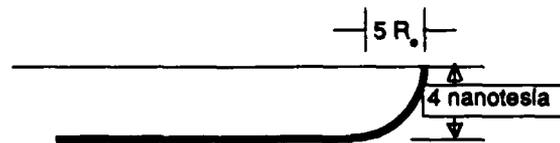
Figure 6 shows the magnetic lines of force within the magnetospheric tail in the noon-midnight meridian plane. The field lines are spaced one degree apart in the high latitude intercept of the earth. They vary in latitude from 60 degrees to 75 degrees. The top panel is the quiet time magnetic field with only the magnetopause, ring and tail currents contributing to the magnetic field. A northward pointing disturbance of 4 nanotesla having a rise time of $5 R_E$ and long persistence is allowed to interact with the magnetosphere. The second panel shows the magnetic field topology when the leading edge of the disturbance pulse has reached the $-20 R_E$. The remaining panels show the changes in the magnetic field topology as the pulse propagates down the tail. The northward disturbance strengthens the z-component of the magnetic field in the tail and thus produces a more dipolar tail field geometry. When a short northward



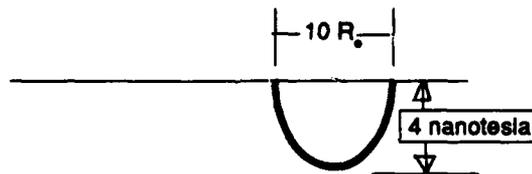
4 nanotesla northward pulse rising to peak value over a distance of $5 R_s$



4 nanotesla northward pulse rising to its peak value in a distance of $5 R_s$ and then decaying again to zero

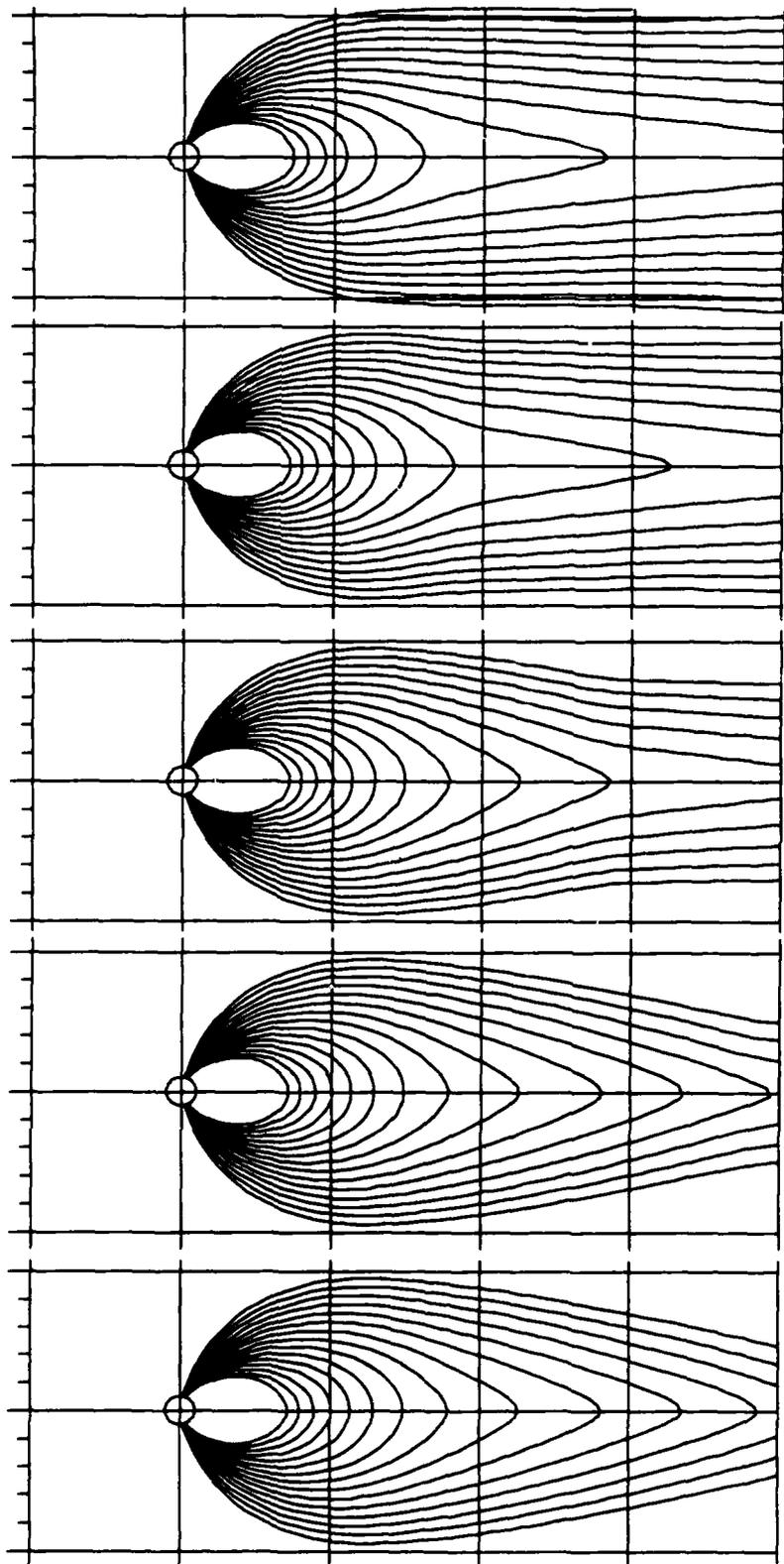


4 nanotesla southward pulse rising to peak value over a distance of $5 R_s$



4 nanotesla northward pulse rising to its peak value in a distance of $5 R_s$ and then decaying again to zero

Figure 5. The above four schematics show the disturbance pulses. The tail magnetic field model was subjected to the four types of disturbances shown above. The pulse moves from left to right at the Alfvén speed.



Quiet time - field lines
using Magnetopause,
Ring and Tail current
systems only

Leading edge of distur-
bance is at $-20 R_E$

Leading edge of distur-
bance is at $-30 R_E$

Leading edge of distur-
bance is at $-40 R_E$

Leading edge of distur-
bance is at $-50 R_E$.
Note: Dipole-like com-
pressed tail.

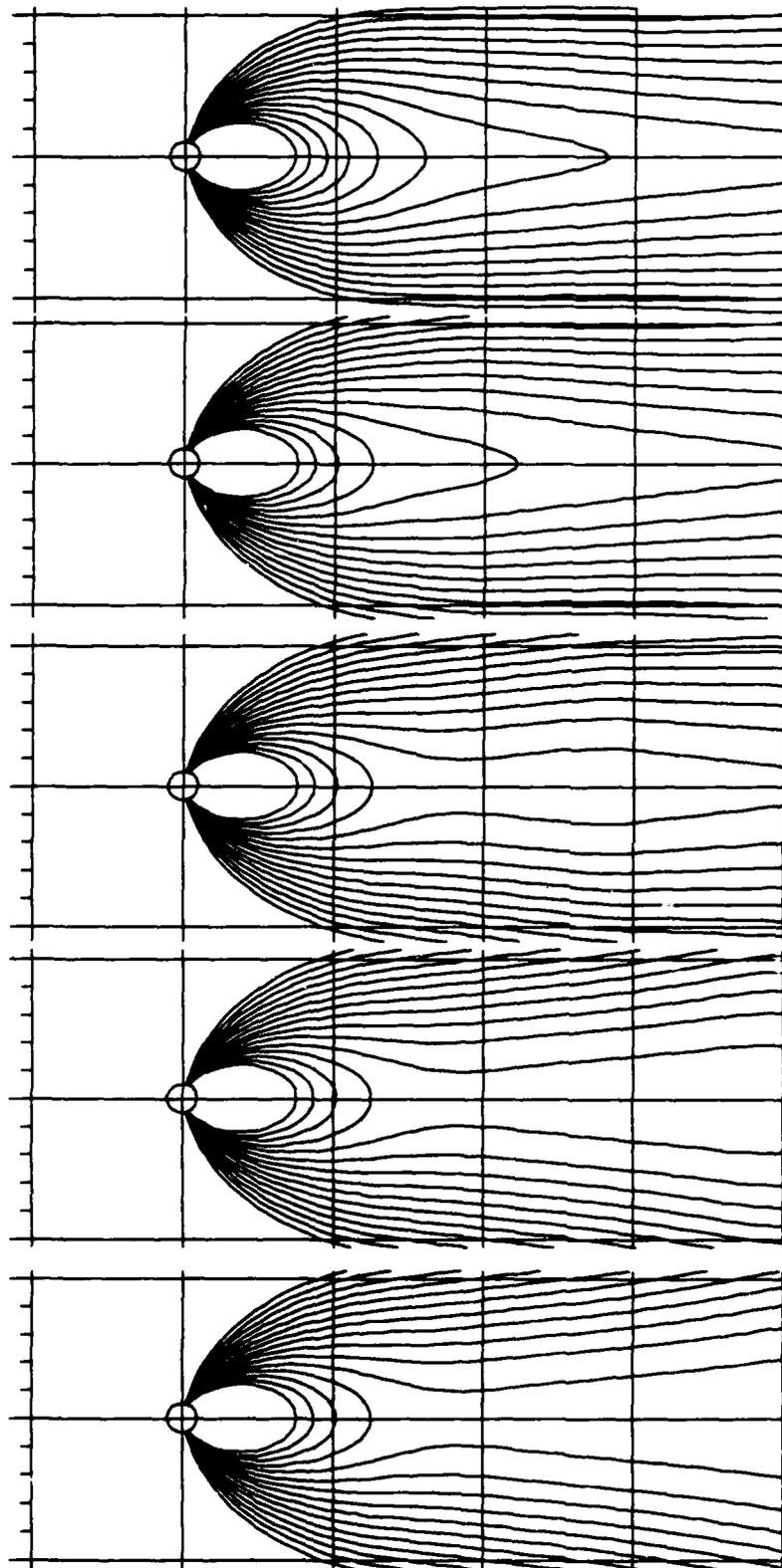
Figure 6. Noon midnight meridian cross section of the tail showing the propagation of a northward pointing disturbance. Pulse is a long 4 nanotesla northward pulse.

disturbance is allowed to propagate down the tail, the results are very similar to the above result with the long pulse. The field becomes more dipolar in a limited region of space. Once the pulse has passed, the field recovers to its initial configuration.

The increasing tail strength due to the entering disturbance wave has many interesting implications. These implications will be studied in more detail during the remainder of this contract. The long pulse reduces the magnetic field gradients along the flanks of the tail and thus reduces the amount of magnetosheath plasma that can enter across the magnetospheric boundary. The enhanced field within the tail increases the magnetic current limit, and thus increases the allowed plasma sheet current densities. Thus the electric field strength in the boundary layer as well as the cross tail electric field must change in response to these variations in the magnetic field and magnetic field gradients. The effect of these changes on the plasma and on the electric fields have not yet been determined.

3.1.2 Southward Pointing Disturbance

Figure 7 shows the magnetic field lines in the noon-night meridian plane when a southward disturbance propagates through the magnetosphere. The disturbance pulse rises over a distance of $5 R_p$ and remains at the disturbance level for an extended time (see figure 5). The top panel represents the quiet magnetosphere before the arrival of the pulse. The next panel shows the tail field topology when the leading edge of the pulse is at $-20 R_p$. Subsequent panels show the leading edge moving in the anti-solar direction. It can be seen that the pulse results in a weakening of the tail and that the tail takes on a more taillike appearance (i.e. the tail becomes extended). When the pulse is at distances larger than $-30 R_p$, the figures show a bulge in the field lines near the neutral sheet (i.e. the field lines no longer appear to cross the neutral sheet). To investigate the behavior of the field lines in this case it is necessary to generate a second set of field line plots, a set where all field lines are forced to originate from the equator. Figure 8 shows the same case but in this figure the displayed field lines originate at the equator instead of the polar cap. One can see that pulse creates an "O" and an "X" type neutral point. There is a set of field lines that is no longer connected to the polar cap and this set of field lines moves in the anti-solar direction along with the disturbance pulse. For this long extended pulse the "O" type neutral point moves down the tail, whereas the "X" type neutral point remains fixed near $-20 R_p$.



**Quiet time - field lines
using Magnetopause,
Ring and Tail current
systems only**

**Leading edge of distur-
bance is at $-20 R_E$**

**Leading edge of distur-
bance is at $-30 R_E$**

**Leading edge of distur-
bance is at $-40 R_E$**

**Leading edge of distur-
bance is at $-50 R_E$.
Note: The bulge in the
tail and that field line no
longer cross the neutral
sheet**

Figure 6. Noon midnight meridian cross section of the tail showing the propagation of a southward pointing disturbance. Pulse is a long 4 nanotesla southward pulse.

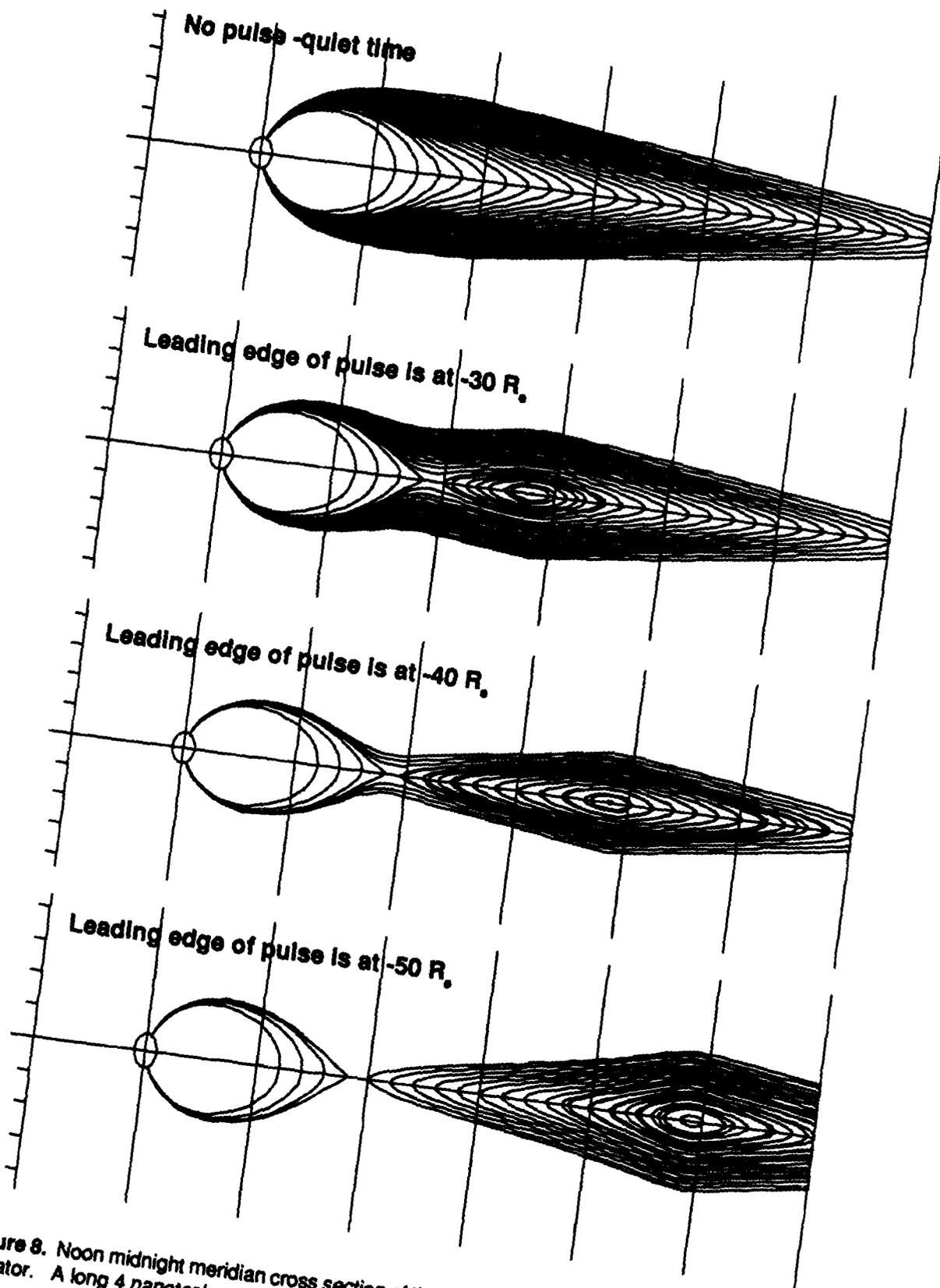


Figure 8. Noon midnight meridian cross section of the magnetosphere. Field lines are drawn from the equator. A long 4 nanotesla pulse moves down the tail. Note the 'X' and 'O' type neutral points.

Figure 9 is similar to Figure 8 except that the disturbance pulse is short. The pulse rises from 0 to 4 nanotesla in $5 R_E$ and then decays to zero over a distance of $5 R_E$. Once again the top panel is the undisturbed configuration and subsequent panels show the disturbance pulse at different locations in the tail. In this figure both the "X" and "O" type neutral point move down the tail at the speed of the disturbance.

The behavior of the field lines when a southward disturbance is introduced is similar to the "plasmoid" topology often described by experimentalists. The southward disturbance significantly weakens the field in the tail neutral sheet. It actually reverses the field over a region of the neutral sheet and thus produces the neutral points. The weakening of the field reduces the amount of plasma the field can support and the magnetically limited currents are drastically reduced. This analysis is a first step and includes only the first order effects of the propagating magnetic wave. It does not as yet include the effects of the plasma response. The plasma response is expected to significantly alter the behavior of the pulse because the beta of the plasma near the neutral point will exceed one unless the plasma rearranges itself in response to the magnetic field.

3.2 Summary of Work During Past Year

The models of the disturbance wave propagating within the tail of the magnetosphere are the first of a series of steps necessary to develop a comprehensive understanding of magnetotail dynamics. This work in conjunction with the work on particle entry into the magnetosphere will provide a detailed understanding of magnetospheric particle energization. The magnetic field model of the disturbance magnetic field when combined with the existing dynamic magnetic field model provides a tool for evaluating the response of the plasma to the magnetic stimulus of the IMF. This effort will ultimately permit us to understand the role of the IMF in controlling and/or triggering processes within the magnetosphere.

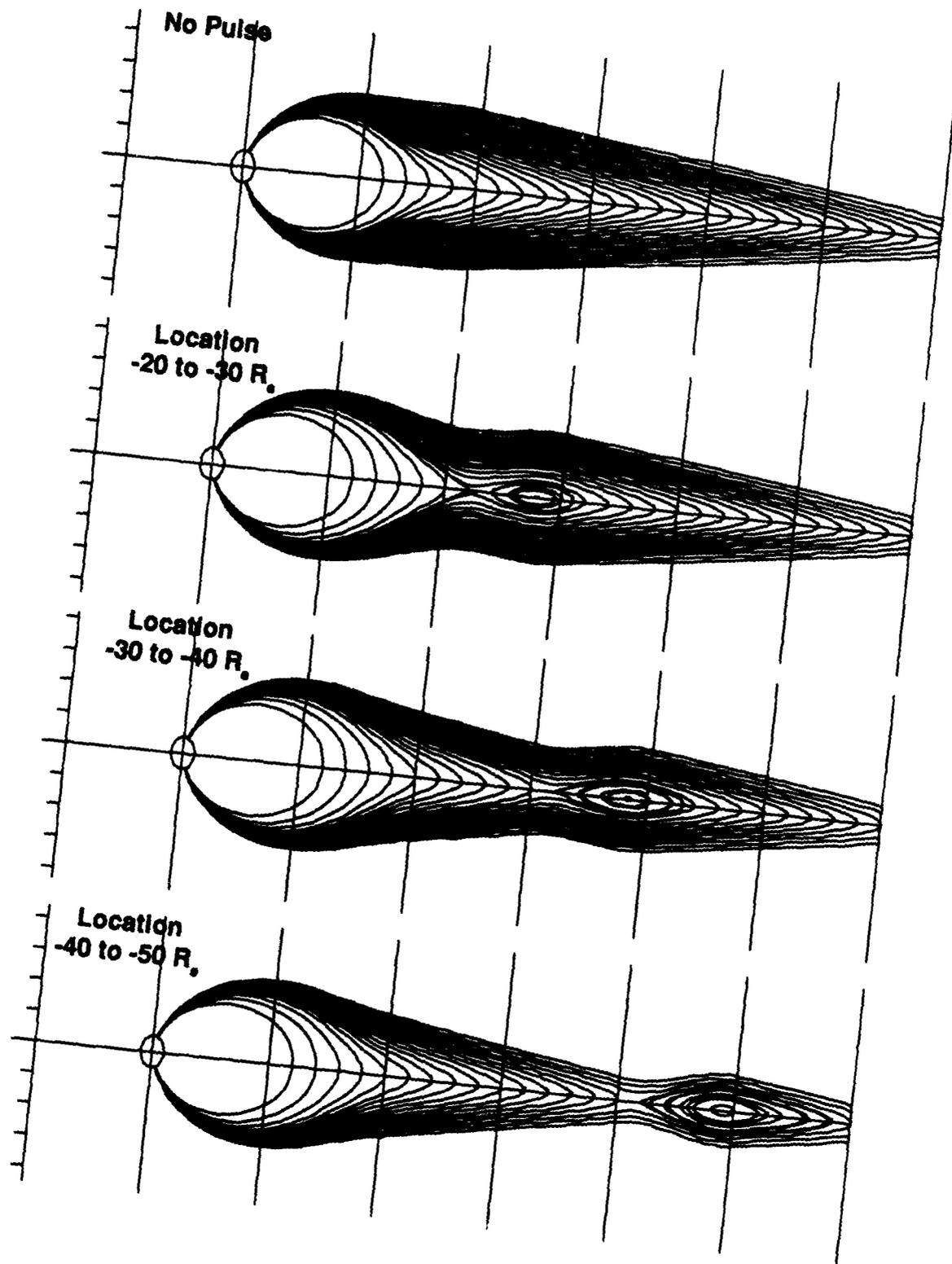


Figure 9. Noon midnight meridian cross section of the magnetosphere. Field lines are drawn from the equator. A short 4 nanotesla pulse moves down the tail. Note the 'X' and 'O' type neutral points. The disturbance pulse rises in a distance of $5 R_e$, and then decays back to zero in a distance of $5 R_e$.

Section 4

PUBLICATIONS AND PRESENTATIONS

Olson, W. P. and K. A. Pfitzer, Use of Geosynchronous Magnetometer Measurements to Estimate Solar Wind Parameters, Presented to the Fall American Geophysical Union Meeting, San Fransisco, CA, December 1987.

Huston, S. L., K. A. Pfitzer, T. Mogstad and W. P. Olson, Prediction of Plasma Parameters in the Low-Latitude Boundary layer and Plasma Sheet Using Gradient Drift Entry, Presented to the Fall American Geophysical Union Meeting, San Fransisco, CA, December 1987.

Pfitzer, K. A., W. P. Olson and T. Mogstad, A Time Dependent, Source Driven Magnetospheric Magnetic Field Model, Presented to the Spring American Geophysical Union Meeting, Baltimore, MD, May 1988.

Olson, W. P. and K. A. Pfitzer, Response of the Magnetotail Magnetic Field Topology to Changes in the North-South Component of the Interplanetary Magnetic Field, Presented to the Spring American Geophysical Union Meeting, Baltimore, MD, May 1988.